



# Laboratory modeling of buoyant jet in rotating fluid

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## **OUTLINE:**

- 1. Introduction to the process
- 2. Laboratory experiments
- 3. Scaling analysis
- 4. Conclusions







1. Introduction to the process

Q: how the Coriolis force and bottom topography influence on the propagation of the surface warm buoyant jet in fresh waters?

Baltic Sea, 24<sup>th</sup> of March, 2000, NASA







#### 2. Laboratory experiments: experimental set-up



1 – rotating table; 2 – slope; 3 – tank wall; 4 – video camera; 5 – autotransformator; 6 – ampermeter; 7 – voltage supply; 8 – copper plate (anode); 9 – heating resistance isolated cable; 10, 12 – outlets; 11 – nichrome wire (cathode).

Experiments were performed: (i) with rotation for rotation periods of 5, 10, 15 s, and Coriolis parameter,  $f = 2\Omega - 0.8$ , 1.25, 2.5 s<sup>-1</sup>, respectively; and (ii) without rotation using various specific power supply q, Wt/m (25.7; 13.1; 6.37 Wt/m) in a tank with slopping and horizontal bottom





**MOVIE 1** 









Case with rotation, slopping bottom, view from above: a) warm buoyant jet colored in dark blue propagating from the wall to the deep central part of the tank; b) velocity field, calculated from paper pellets displacement using Tank Field Calculator program. Time from the start of the run - 35 min; specific power supply q = 25.7 Wt/m.





**MOVIE 2** 







Case with rotation, horizontal bottom, view from above: a) warm meandering temperature front at q=25.7 Wt/m and Coriolis parameter  $f = 0.8 \text{ s}^{-1}$ ; b) calculated surface velocity field with baroclinic eddies marked by ovals.





#### ? Radial velocity of the jet propagation, U,m/s ?

#### Main external dimensional parameters:

- buoyancy flux (initiated by line source), B,  $m^3/s^3$ ;
- Coriolis parameter, f, s<sup>-1</sup>;
- depth of the tank, H, m;
- the time, t, s;
- thermal diffusivity,  $k_T$ ,  $m^2/s$ ;
- kinematical viscosity, v,  $m^2/s$ .





Results of laboratory experiments have shown that:

 $U \sim B^{1/2};$ 

 $U \sim f^{-1};$ 

 $k_{\rm T}\,$  and  $\nu$  are nearly constant;

U does not depend on time.

Main dimensionless parameters governing the process:

$$Ra_F = \frac{g\alpha FH^3}{\rho_0 c_p \kappa_T^2 v}$$

- (flux) Rayleigh number or

$$Ra_B = \frac{BH^3}{\kappa_T^2 \nu}$$

$$Ek = \frac{V}{fH^2}$$

- Ekman number





Therefore, we may parameterize the velocity of the buoyant jet as the follows:

$$\frac{U}{B^{1/3}} \sim \left( Ra^{1/2} \cdot Ek \right)$$

In the final form, the formula for the radial velocity of the jet is defined as:

$$\frac{U}{B^{1/3}} \sim \frac{B^{1/2} v^{1/2}}{k_T H^{1/2} f}$$

or, using Pr number,

$$U \sim \left(\frac{B^{1/3}}{k_T f^2 H}\right)^{1/2} \Pr^{1/2}$$

where Pr number is constant and for the experiments Pr~7







Velocity of radial propagation of the warm buoyant jet, U, is non-dimensionalized by versus Rayleigh and Ekman numbers in double logarithmic scale. Data of the described experiments in presence of the slope configuration for rotation rates f = 2.5; 1.25; 0.8 s<sup>-1</sup>; correlation coefficient R<sup>2</sup>=0.7488.





Results of laboratory experiments for non- rotating case have shown that:

 $U \sim B^{1/2};$ 

 $k_T$  and v are nearly constant;

U does not depend on time.

#### Main dimensionless parameter:

$$Ra_F = \frac{g\alpha FH^3}{\rho_0 c_p \kappa_T^2 v}$$

- (flux) Rayleigh number or

$$Ra_B = \frac{BH^3}{\kappa_T^2 \nu}$$





By analogy with rotating case, we should parameterize the velocity of the buoyant jet as the follows:

$$\frac{U}{B^{1/3}} \sim Ra^{1/2}$$

Thus, the formula for the velocity jet is obtained in final form:

$$U \sim \left(\frac{B^{1/3}H^3}{k_T^2 \nu}\right)^{1/2}$$







Non-dimensional radial velocity of the warm buoyant jet is as a function Rayleigh number in the double logarithmic scale for non-rotating case. Non-dimensional radial velocity of jet propagation depends on Rayleigh number by the law of  $U_{B^{1/3}} = 4 \cdot 10^{-8} Ra^{1/2}$  with correlation coefficient R<sup>2</sup>=0.767





## 4. Conclusions

- 1. The rotation of the system strongly affects the characteristics of the warm buoyant jet and the velocity of its radial propagation. It has been revealed, that in rotating fluid a strong along-wall cyclonic current is formed. The radial velocity of buoyant jet propagation is about an order of magnitude less than in non-rotating fluid.
- 2. The bottom slope stabilizes the propagation of the temperature front of the buoyant jet, preventing for its breaking-up and formation of the baroclinic eddies. Upon the horizontal bottom radial propagation of the buoyant jet is effected and accelerated by the lateral eddy diffusivity.





## 4. Conclusions

3. Scaling analysis has show that Rayleigh and Ekman numbers are the key non-dimensional parameters that determines the regimes and regularities of the warm buoyant jet in the rotating case, but only Rayleigh number is important for the non-rotating case. The processing of the experimental data revealed that the non-dimensional radial velocity of jet propagation depends on Rayleigh and Ekman numbers by the law of  $U_{B^{1/3}} = 3 \cdot 10^{-4} Ra^{1/2} \cdot Ek$  for rotating case and by the law of  $U_{R^{1/3}} = 4 \cdot 10^{-8} Ra^{1/2}$  for non-rotating case.





## Thank You for attention!!!